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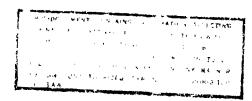
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US 4200 2000 REV. 0/6

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APPLIED PHYSICS RESEARCH

Rhodes and Blorsom

7343 DERING AVENUE
CANDGA PARK, CALIFORNIA
DIAMOND 0.2707

Feb. 1, 1963

FINAL REPORT

DYNASOAR TESTING FOR THE BOEING COMPANY
UNDER P.O. B-438883-9155
WORK STATEMENT 2-5781-4-129



ABSTR4CT:

This report with the models submitted constitutes partial performance under

DYNASOAR TESTING FOR THE BOEING COMPANY UNDER P. O. 8-438883-9155 FINAL REPORT

WORK STATEMENT 2-5781-4-129

ABSTR:CT:

This report with the models submitted constitutes partial performance under the above mentioned P.O. and work statement.

computations and calibrations for the two run configurations, 2) duals with the test log, This report consists of three parts: 1) deals with the details of data reduction, flow installation drawings, model photographs and heat transfer distributions obtained, 3) deals with the theory of the paint operation.

APPROVED:

Willowson, y Project Engineer Rhodes and Bloxsom D. Blox som, Jr.

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02-80911

CONTENTAL

DYNASCAR TESTING

PART I

The following runs were made under the conditions listed:

a. 55° sweep; 5°, right 20°; s zem yaw left	., 20° right; : .erc yaw left	adius, 60° sweets right 20°; stro yow left,	Swaep; Syright 20°; zao yaw left,
Fin Canfiguration 6" diameter leading edgin. 55° sweep; nudder deflection: left 15°, right 20°; yaw of 3.5° right producer zern yaw left 7.0° yinght on model	6" diamater, 55° sycep; rudder reflection, 13° lett, 20° right; yaw of 0,5° right producer sere yaw left 7.0° right on model	8" diameter leading edge rasius, 60° sweep, rudder deflection; left 10°, sight 20°; you of 3.5° right procluces zero yow left, 7.0° yow right on model	6" diameter leading edge, 5.° swrep; rudder deflection 100 left 13", right 20°, yrw of 3.5° right producer zero yaw left, 7.°C yaw right on model
dynes/cm 97000	97000	97000	25,030 6
configuration (14 scale	. 04 scale A	.04 scole A	.04 scale 8
Ann angle of attack configuration 1/2 events	20%	500	20%
grang i	6	ra	×

were andlyzed by means of measured pressure rise in the stagnation chamber. Inte date shown in tables for equilibrium by using a 99% efficiency factor of energy conversion from capacitors to air. The validity of the latter assumption may be seen in Figure 1, in which 38 previous runs known it an temperature and initial bomb pressure (.5%) set up by heurdon and marsh gouges calibraid by dead weight, the stagnation conditions can be computed from air real gas Figure 1 her a scatter of 1%, earlier data and lyzed for 59 additional runs shown scatter The stagnation conditions were set up by measurement of the coaccitor vallage to .5% and the capacitance to 1%. With the known bomb volume (.5%) and the between 75 and 125% officiencies.

obtained from previous experience as to nozzle velocity and dynamic pressure. These cullibrations are encised. Dynamic pressure was measured on Run 4. The rest of the nozzle parameters Once the stagnation conditions were known, then the nozzle collibrations were



Yaw of 0.5° right producer Lerc yaw left 7.0° right on model

8" diameter leading edge tatius, 60° sweeps redder deflection: left 10°, sight 20°; you of 3.5° right produces zero you left, 7.0° yaw right on model

.04 scale

8

7.0% yaw right on model
.04 scale 25,000 6" diameter leading coge, 5% swaep;
.04 scale 25,000 6" diameter leading coge, 5% swaep;
.v.cdc.deficction 100 left 10%, right 20%,
yrw of 3.5% right produces zaro yaw left,
7.0% yaw right on model

500

were and fixed by means of ansaured pressure rise in the stagnation chamber. Tate date shown in tables for equilibrium by using a 99% efficiency factor of energy eurversion from expression to air. The validity of the latter assumption may be seen in Figure 1, in which 38 previous rous known it an temperature and initial komb pressure (.5%) set up by heurdon and marsh youges calibrated by dead weight, the stagnation conditions can be composed from cirrieal gos Figure 1 hat a scatter of 1%, earlier data analyzed for 59 additional runs shown scatter The stagnation conditions were set up by measurement of the copacitor voltage to .5% and the capacitance to 1%. With the known bomb volume (.5%) and the between 75 and 125% afficiencies.

obtained from previous experience as to nozzle velocity and dynamic pressure. These cullibrations are encicsed. Dynamic pressure was measured on Run 4. The rest of the nazzle parameters Once the stagnation conditions were known, then the nozzla calibrations were were obtained from computation from the knewn frozen flow nozzle parameters.

the spherrs. Spheres of various diameters from 3.00 inches to .250 inches were used to determine equilibrium sphere flow of . 070R (where R is sphere radius).* Hence the flow is in equilibrium on page 259, vol. 26 (4) was used to compute the stagnation retas on the ed ibration spheras, and was . 063 R checking the expected shock standoff distance for non-equilibrium rozzle flow and for shock thindoff distance to determine flow equilibrium conditions on the calibration sphere. the heat transfer rates to be assigned to the varietus colons. Tal readings of these ratios of heat transfer are plotted in the color curve of part 3). Once these rates were obtained, the same heat transfer contours were alt referenced to the heat transfee rates expected on the Dynasoar the calibrating spheres. Onco this was known the theory of Lees (Jet Propulsion, April 1956 the same theory with measured pressure gradient (figure 2) was used to compute rates around colors were read on the models and surface heat transfer contours developed. These surface From previous calibrations, the shock standoff distance measured for 5 spheres in 5 runs Once the nozzle parameters were known, the calibration sphere was checked alider at the following conditions:

altitude 250, 000 feet (ARDC 1959 model atmosphere) velocity: 20, 000 feet/second

All the developed plots refer to these conditions, but actual rates in the tunnel may be astained from the scaling factors plotted in configurations A, B.

Due to the short time models were available for reading at the Rhodes and Bloxsom plant, only maximum and calibrations rates were measured and included in this report.

See figure

CONFERNIAL



DYNASOAR TESTING

CALIBRATIONS FOR CONFIGURATION A

 Z_{00}^{∞} constant at 1.32, effective gamma frozen flow: 1.455 density variations given in Figure 4 (serie:) Test section distributions: Voo constant at 15,200 feet/second

temperature variations vary as square root of density variations (similar to

heat transfer variations) To 800GoK, H $_{0.2}$ x108 ft²/sec. ²; P $_{0.2}$ 1405 psia; M $_{0.2}$ 15.2; T $_{0.2}$ 140°K, V $_{0.0}$ 15,200 ft/second tho $_{\infty}$ 8.62 × 10⁻⁷ gm/cm 3 ; $_{2\infty}$ 6.85 × 10⁻³ psiG; Re $_{0o}$ /ft. 124,009; $_{2\infty}$ 1.32 where $_{
m c}$ conditions are stagnation, $_{
m co}$ anditions are tunnel ambient free stream at tunnal nozzle CL

theory: dynamic pressure x 2 equals sphere pressure of .185 atm., density behind sphere shock is 7.6 x 10⁻⁶ cm. cm³. Lees thoory gives 367 watts/cm² for the stagnation point of a 3.00 inch sphere or 523 watts/cm² for the stagnation point of a 1.48 inch sphere (ping pong ball). For a Ferri number of 8630 we find an increase vorticity factor of 1.09 for a final rate of 400 watts/cm² for the 3.00 inch sphere or 570 watts/cm² for the 1.48 inch sphere. The following rates with vorticity Additional parameters computed behind normal shock for equilibrium flow heat transfer using Lee's hold for the flight rates on the scaled 1.48" clameter/.04 sphere:

Flight rates scaling factor 31.7 watts/cm² .04

CALIBRATIONS FOR CONFIGURATION B

CALIBRATIONS FOR CONFIGURATION A

Test section distributions: V_{oo} constant at 15,200 feet/second
2,00 constant at 1.32, effective gamma frozen flow: 1,455 density variations given in Figure 4 (series)

temperature varictions vary as square root of density variations (similar to

heat transfer variations) To 80000 K, H_o 2 ×10⁸ ft²/sec. ²; P_o 1405 psia; M_{oo} 15.2; T_{oo} 140⁶K, V_{oo} 15,200 ft/second

 $^3 \text{ Po}_{\infty} 8.62 \times 10^{-7} \text{ gm/cm}^3; p_{\infty} 6.85 \times 10^{-3} \text{ psia; Re}_{\infty}/\text{ft. } 124,009; Z_{\infty} 1.32$

where c conditions are stagnation, where connected the stream at tunnal nozzle CL

theory: dynamic pressure x 2 equals sphere pressure of .185 atm., density behind which shock is 7.6 x 10⁻⁶ cm. cm³. Lees theory gives 367 watts/cm² for the stagnation point of a 3.00 inch sphere or 523 watts/cm² for the stagnation point of a 1.48 firch sphere (ping pong ball). For a Ferri number of 8630 we find an increase vorticity factor of 1.09 for a final rate of 400 watts/cm² for the 3.00 inch sphere or 570 watts/cm² for the 1.48 inch sphere. The following rates with vorticity Additional parameters computed behind normal shock for equilibrium flow heat transfer using Lee's hold for the flight rates on the scaled 1.48" chameter/.04 sphere:

Flight rates scaling factor 31.7 watts/cm2

CALIBRATIONS FOR CONFIGURATION B

Test section distributions: V_{oo} constant at 14,500 feet/second Z_{∞} constant at 1.36, effective gamma frozen flow: 1.495 density variations given in Figure 5

temperature variations similar to Figure 4 heat transfer variations T_{σ} 750C°K, P_{o} 350 psia, H_{o} 2×108 ft²/sec²; M_{o} 15.4, T_{o} 138°K, V_{oo} 14,500 ft/sac

tho. 2.55 x 10-7 gm/cm³, p.o. li.32 x10-4 otm. ; Re /ft. 31,000; Zoo 1.32

Referring the rates to the flight: conditions with vorticity we get for the scaled 1.48" Flight rates scaling factor 31,7 watts/cm² .04 diameter/.04 sphere:

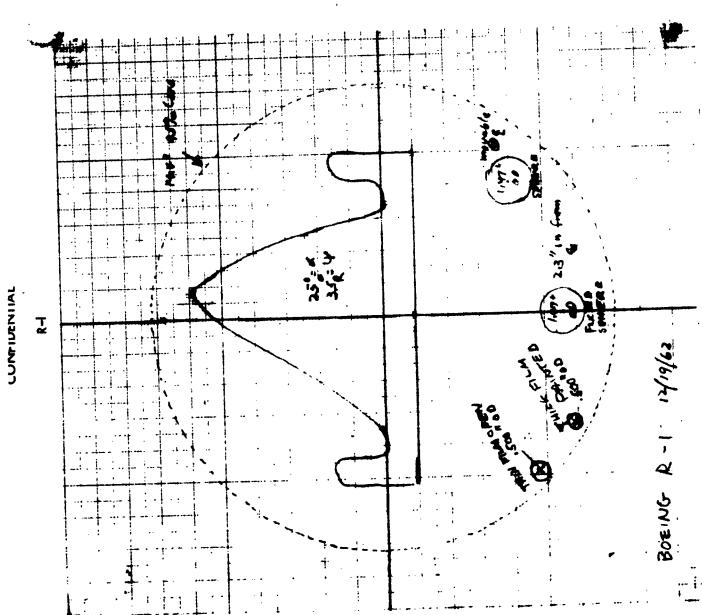
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5500 frames/second movies are also included of side view in self luminosity light.

Figure 6 gives the theory and calibration of the paint operation (Part 3).









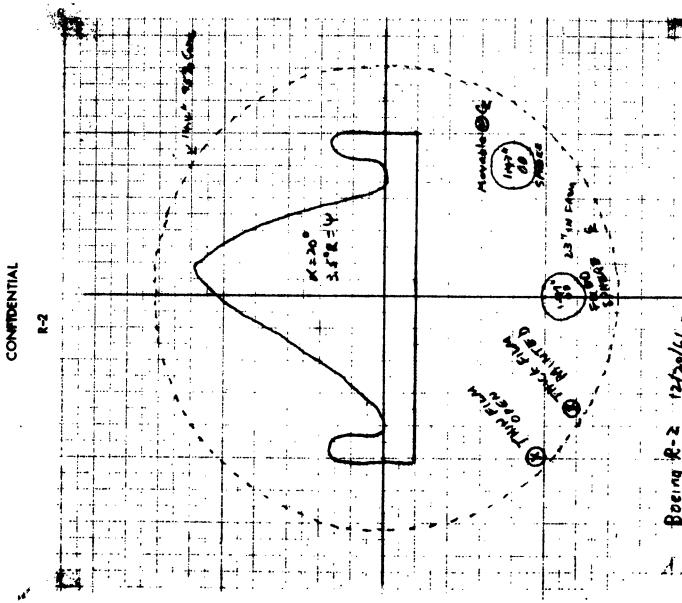


INSTALLATION DRAWING

BOEING R-1 1419/63

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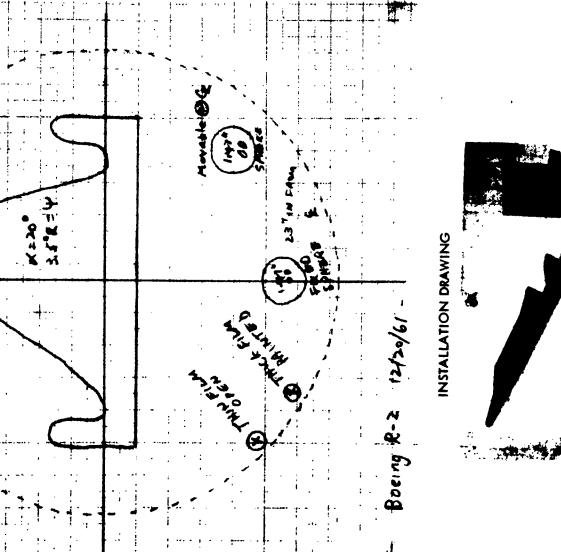
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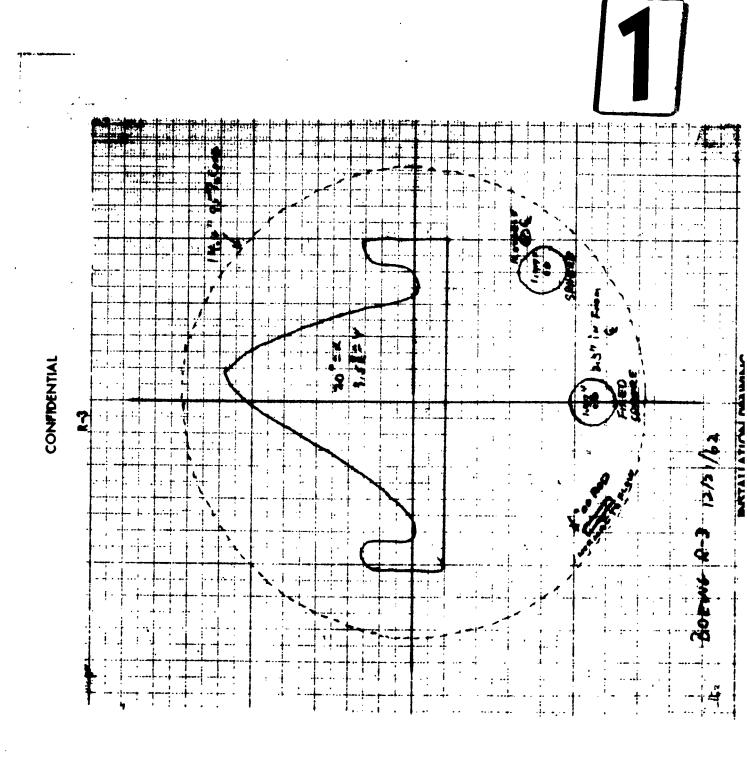






INDOEL PHOTOGRAPH

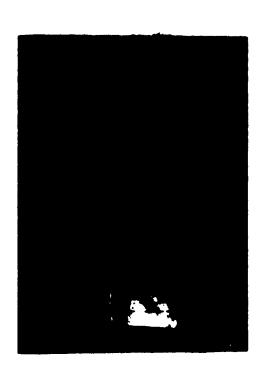








MODEL PHOTOGRAPH



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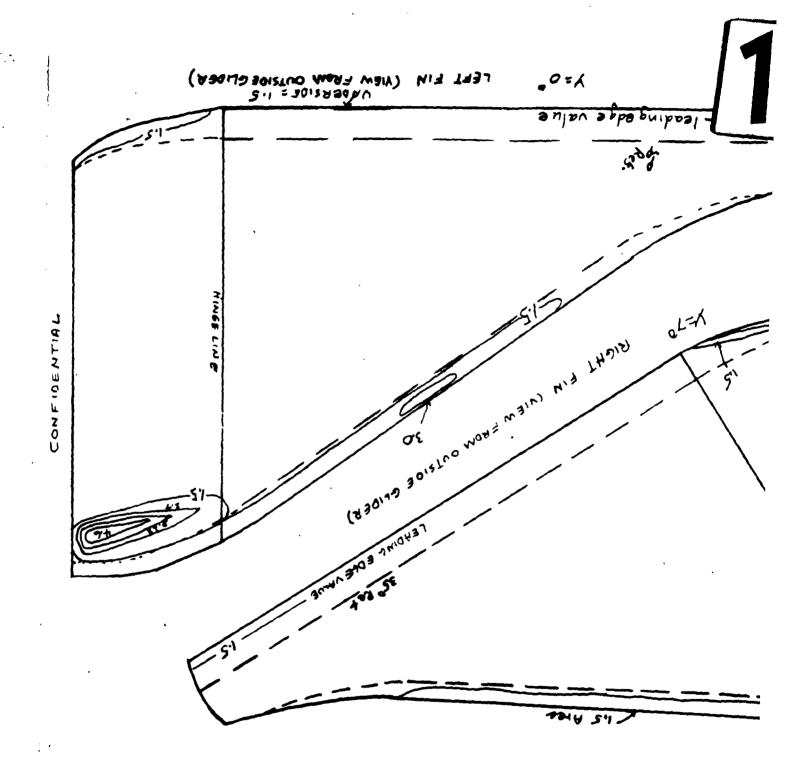
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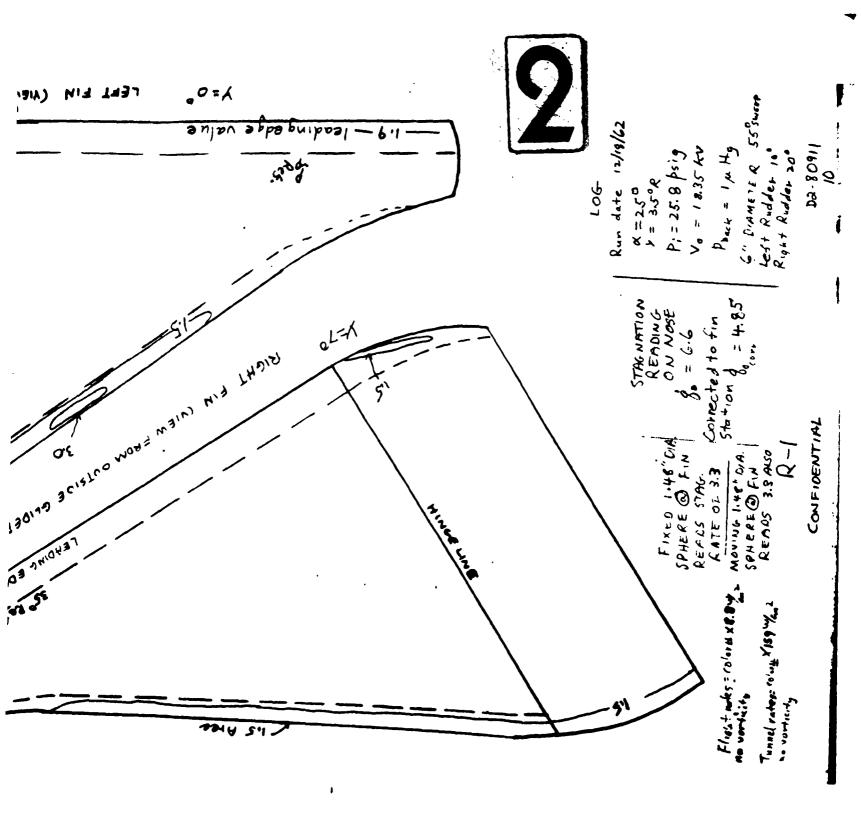
Boeing R-4 1221/62

20°=2 3.5°= 4.

MODEL PHOTOGRAPH

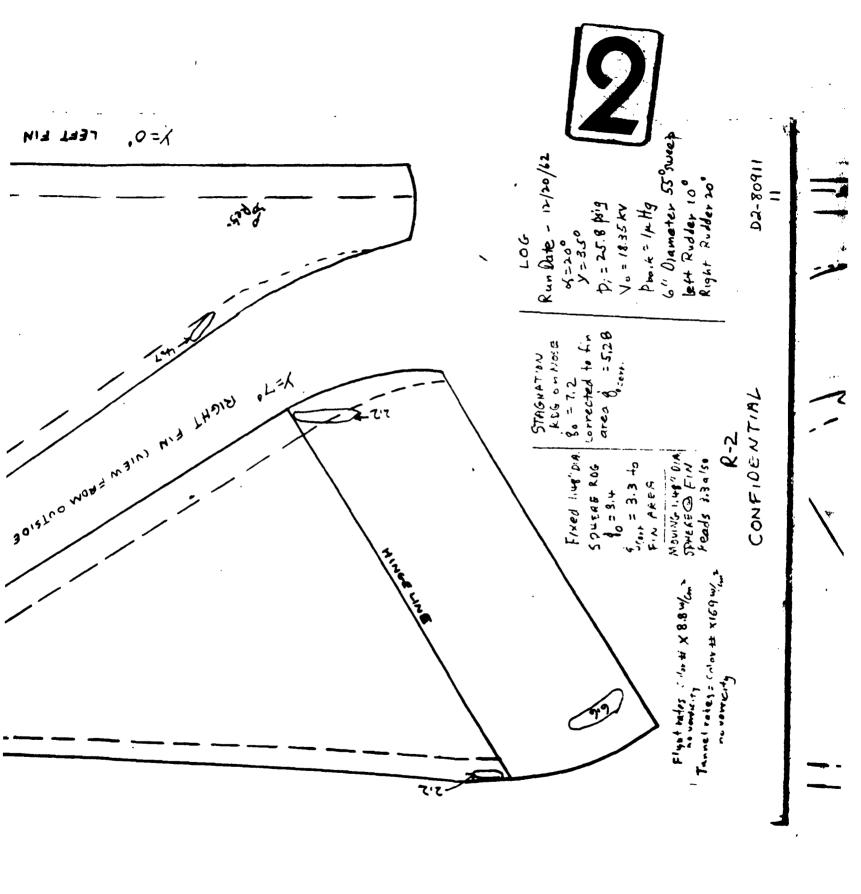
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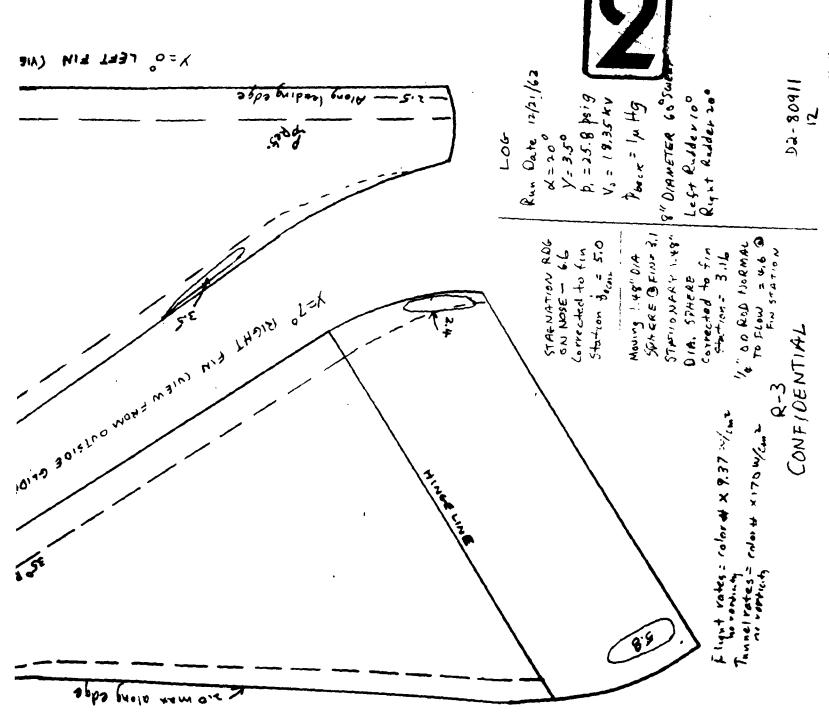
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TEET FIN (VIEW FROM OUTSIDE GLIDER)

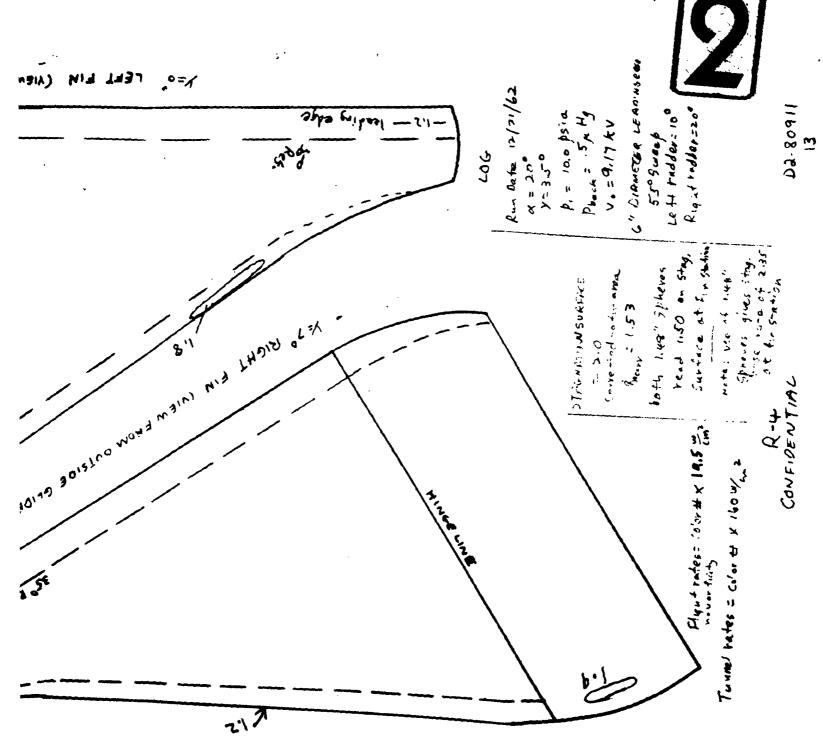
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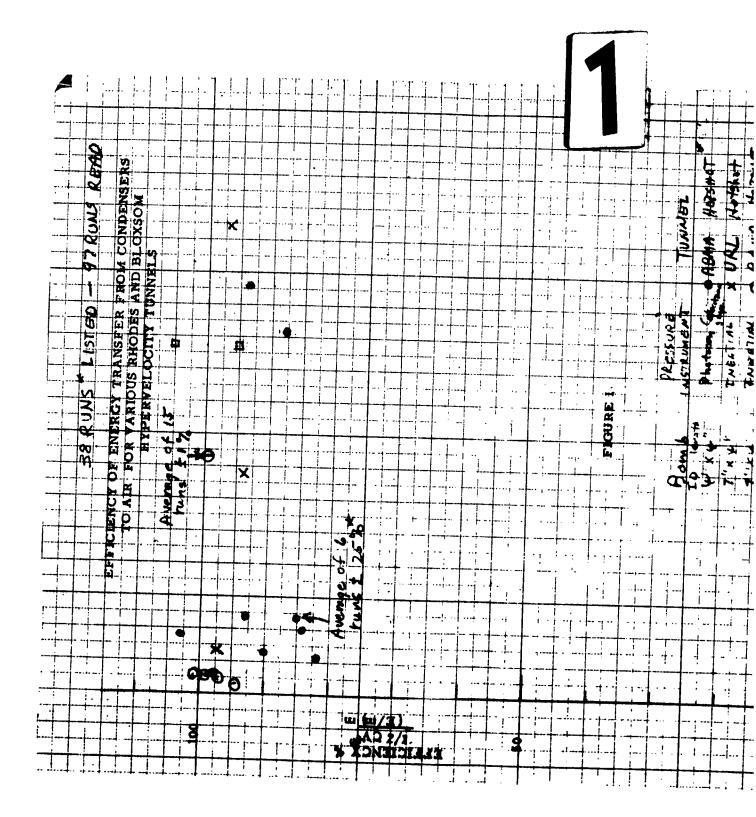


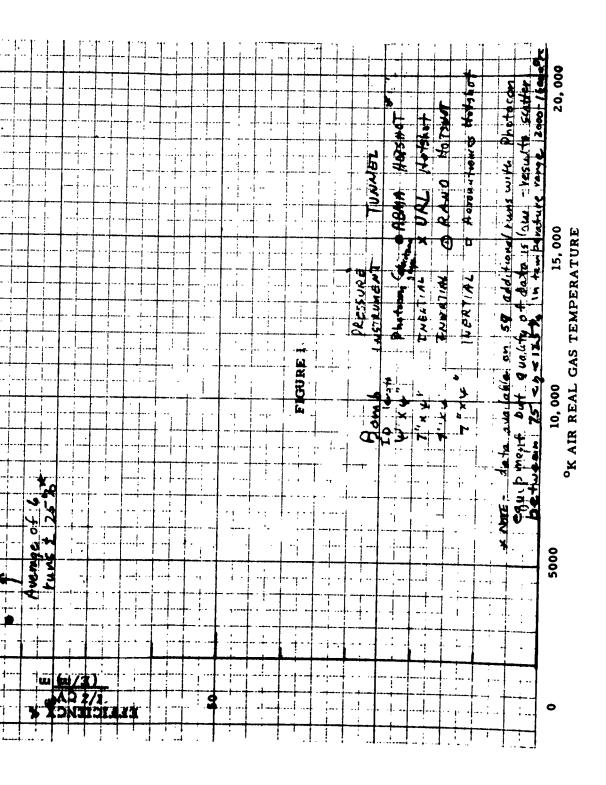
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LEFT FIN (YIEW FRAN OUTSIDE GLIDER)

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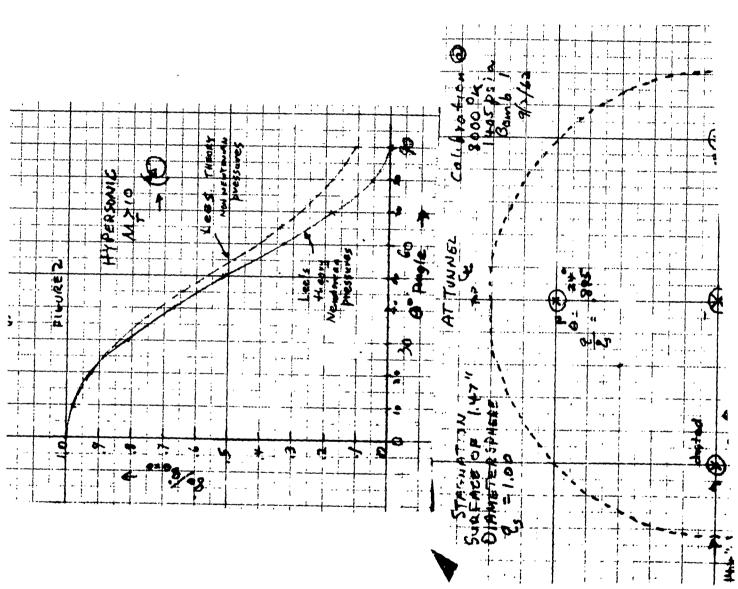




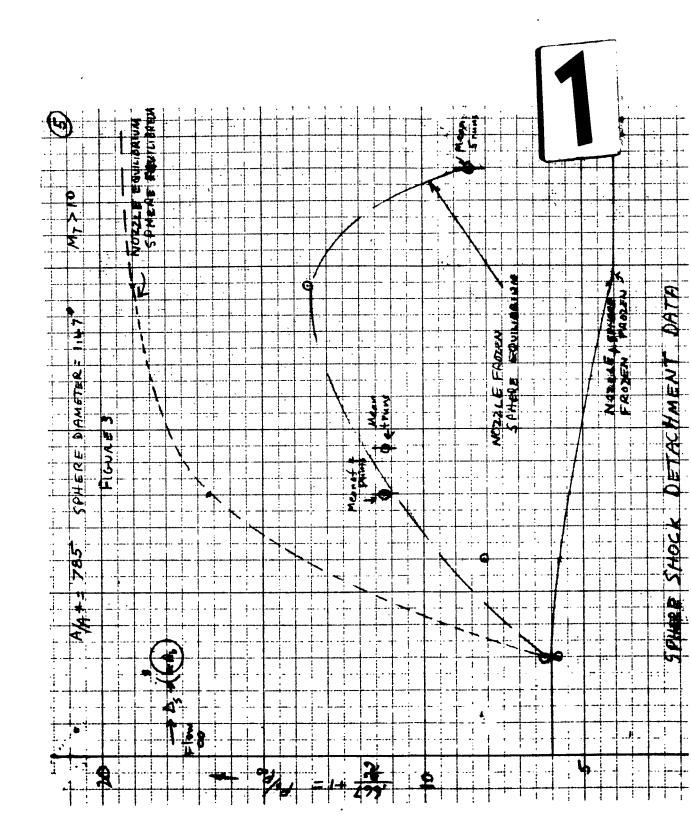


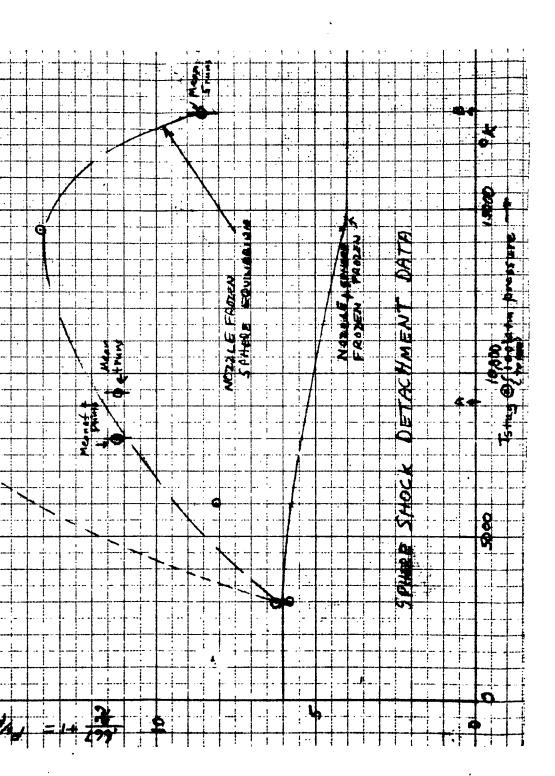














Rhodes and Bloxsom

APPLIED PHYSICS RESEARCH

7343 DEERING AVENUE CANOGA PARK. CALIFORNIA

DIAMOND 0.2707

FIGURE 4 SERIES

January 4, 1963

CALIBRATION SUMMARY: 20.5" NOZZLE, 72" STATION . 750" THROAT

I) ISENTROPIC CORE: 3 RUNS (HEAT TRANSFER)
1 RUN (DYNAMIC PRESSURE)

II) LONGITUDINAL CALIBRATION: 2 RUNS (HEAT TRANSFER) 1 RUN (DYNAMIC PRESSURE)

III) FLOW ANGULARITY: 2 RUNS



CALIBRATION SUMMARY: 20.5" NOZZLE, 72" STATION . 750" THROAT

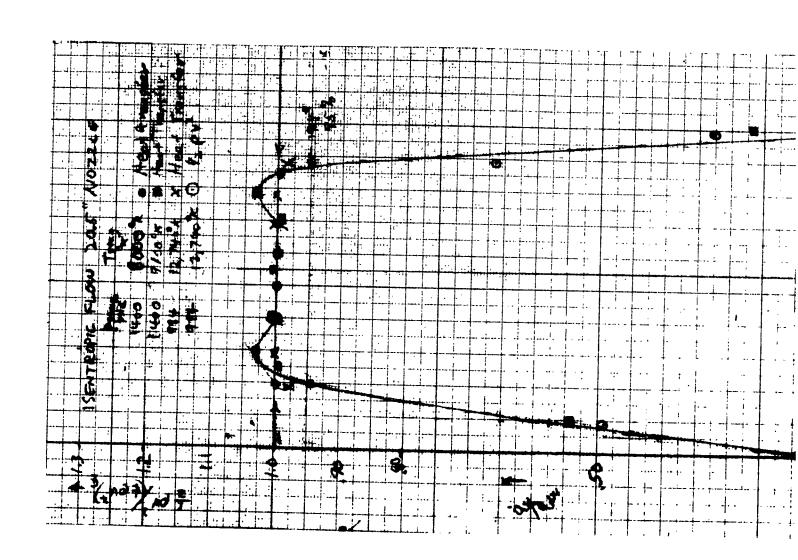
I) ISENTROPIC CORE: 3 RUNS (HEAT TRANSFER)
1 RUN (DYNAMIC PRESSURE)

II) LONGITUDINAL CALIBRATION: 2 RUNS (HEAT TRANSFER)
1 RUN (DYNAMIC PRESSURE)

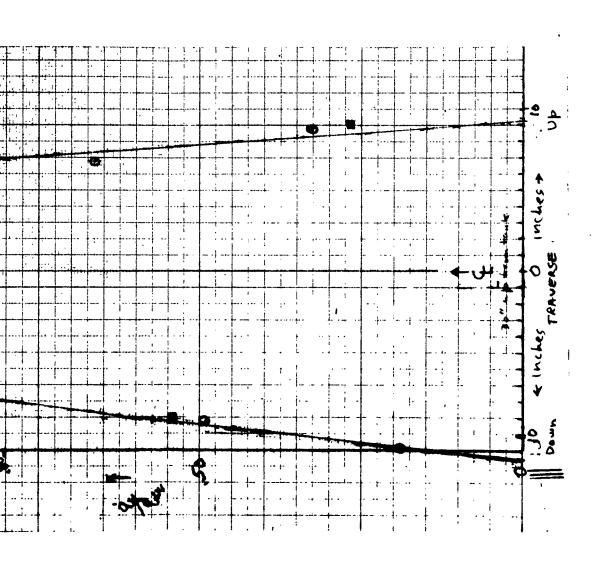
III) FLOW ANGULARITY: 2 RUNS

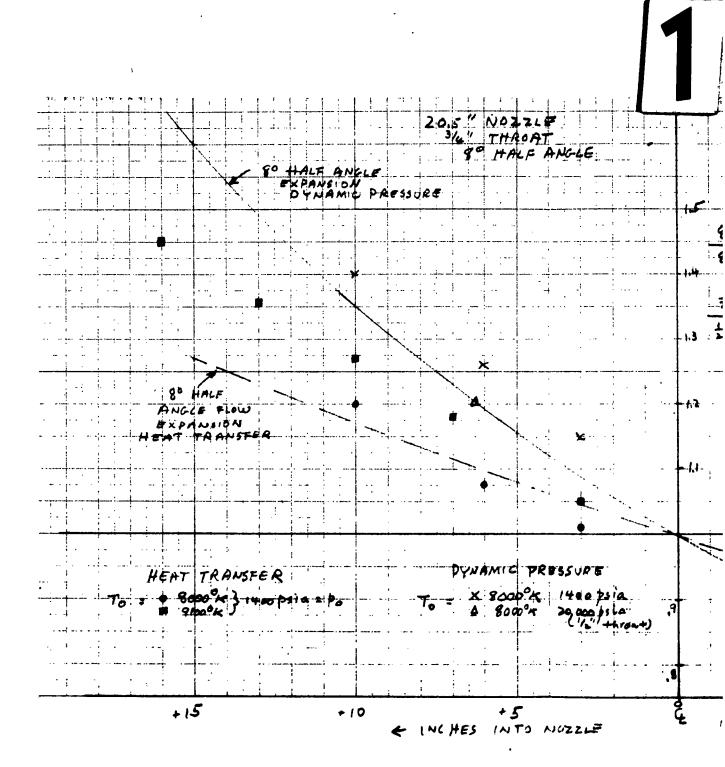


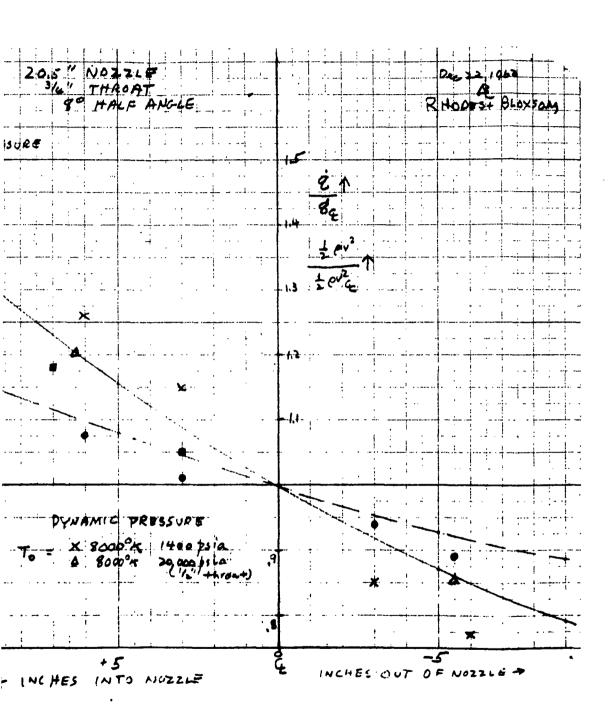
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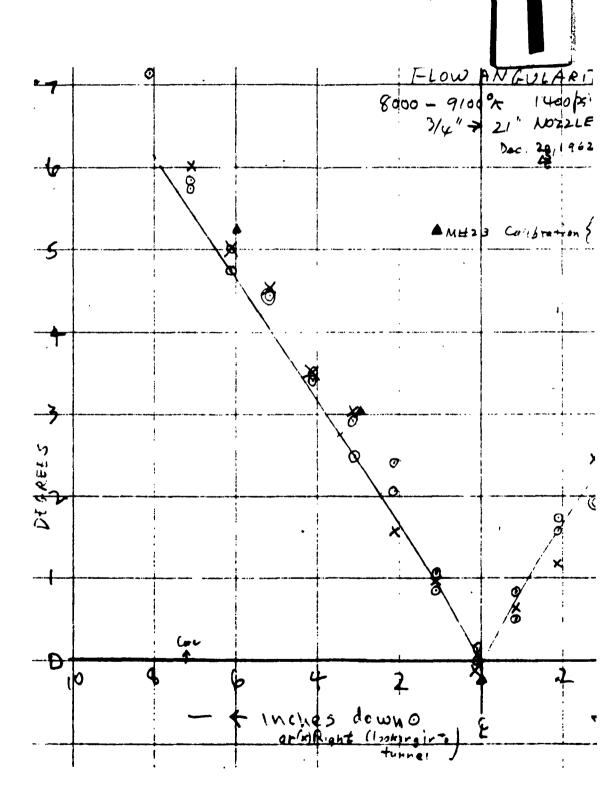


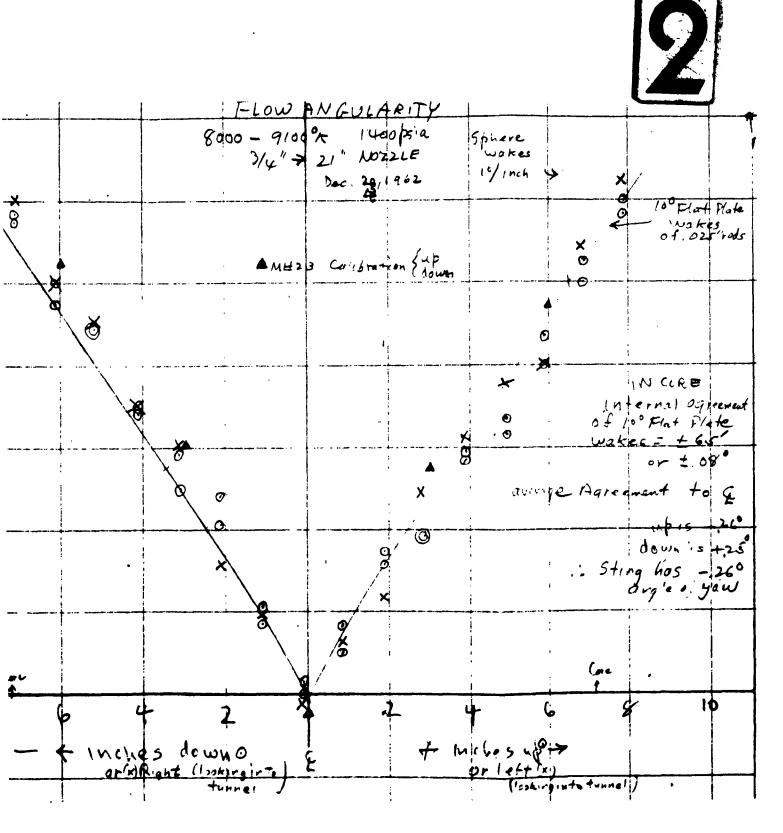












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Rhodes and Bloxsom

APPLIED PHYSICS RESEARCH

7343 DEERING AVENUE CANOGA PARK, CALIFORNIA Білмомь 0-2707 January 4, 1962

.

CALIBRATION B Figure 5

CALIBRATION REPORT

20. 5 INCH NOZZLE, 72 INCH STATION, . 750 INCH THROAT

and 5. 5 INCHES TO RIGHT AND LEFT AT 72 INCH STATION I) BENTROPIC CORE: 1 RUN (HEAT TRANSFER) UP AND DOWN CENTERLINE

II) LONGITUDIMIL CALIBRATION: (not measured, probably same as 1400 ps; III) FLOW ANGULARITY: calibrations)

Reynolds number/ft 31,000

Amblent density 255 x 10-7 gm/cm3 (measured from sphere acceleration dynamic pressure)

H 2 x 10 4 ft 2/sec 2

stagnation pressure 350 psia

stagnation temperature 7500°K.

1

CALIBRATION REPORT

20. 5 INCH NOZZLE, 72 INCH STATION, . 750 INCH THROAT

I) BENTROPIC CORE: 1 RUN (HEAT TRANSFER) UP AND DOWN CENTERLINE and 5. 5 INCHES TO RIGHT AND LEFT AT 72 INCH STATION

(not measured, probably same as 1400 psi calibrations) II) LONGITUDIMIL CALIBRATION: III) FLOW ANGULARITY:

Reynolds number/ft 31,000

Amblent density 255 x 10-7 gm/cm3 (measured from sphere acceleration dynamic pressure)

H 2 x 108 ft 2/sec 2

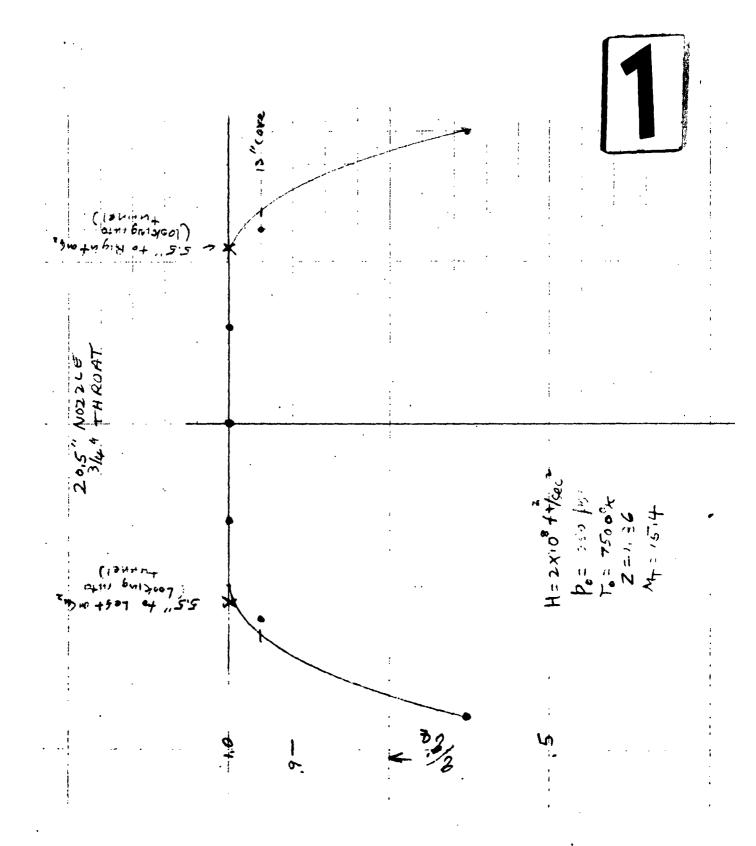
stagnation pressure 350 psia

stagnation temperature 7500°K.

Z 1. 36

Mach number 15.4





,3

H=2x108 ++/cec 10= 300 from 10= 7500 % Z=1,36 Inches traverse up >

STANDO

MEDIFIC & ESSER CO.

THEORY OF PAINT OPERATION

films, such as oil slicks on water, and are due to selective interference and These colors are identical in hue and intensity with those obtained from thin Readout of heat transfer rates are accomplished by means of a series of colors which vary from violet to red through two or three separate changes. reinforcement of various wayelengths present in the white light readout.

three hundred years and is reported in various physics books (Physics, Hausmann, The theory of operation of these colors from thin films has been known for Slack, Second Edition, pages 679-685).

X = nx . n=1,3,5... The series of reinforcement colors is given by

the series of interference colors is given by (normal incidence)*

X-12/020 X-14,... where x is the film thickness in Angstrom units (10^{-8} cm.), and μ is the index of refraction of the film and λ is the light wavelength in Angstrom units. (normal incidence)*

For a black surface, the n equals one results in no light reflected again due For a black surface, the n equals zero results in no light reflected. to the lack of interference with the colors not reinforced.

reference is the n equals 2 series of interference colors. The first of these colors, The thinnest film that will appear colored, as explained on page 683 of above corresponding to an interference of violet light is a reddish color produced by the combination of all other colors. This red is not presently picked up due to its faintness, these colors are present in inverse order.

The n equals 3 order of colors is very bright, as the n of 3 order is reinforcing while the n equals 4 order is supresseing other colors, and the calibration of this order is seen in the accompanying figure for the equation: hin Angstromanits

x = nx film the knoss

= 1.30 ×10 (H) - 3.0

three hundred years and is reported in various physics books (Physics, Hausmann, The theory of operation of these colors from thin films has been known for Slack, Second Edition, pages 679-685).

TIPE TEAMORE

The series of reinforcement colors is given by $X = n\lambda \cdot n = 1.3.5$.

(normal incidence)*

the series of interference colors is given by $X = A \lambda = 0, 2, 4, \dots$ (normal incidence)*

(normal incidence)* where x is the film thickness in Angstrom units (10^{-8} cm.), and μ is the index of refraction of the film and λ is the light wavelength in Angstrom units. (normal incidence)*

For a black surface, the n equals one results in no light reflected again due For a black surface, the n equals zero results in no light reflected. to the lack of interference with the colors not reinforced.

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The n equals 3 order of colors is very bright, as the n of 3 order is reinforcing while the n equals 4 order is supresseding other colors, and the calibration of this order is seen in the accompanying figure for the equation:

1)
$$\frac{2}{\xi_{i}} = 1.30 \times 10^{-3} \frac{11 \times 1}{4 \times 1} - 3.0 \times 10^{-3} \frac{11 \times 10^{-3} \times 10^{-3}}{11 \times 10^{-3}}$$

thecelor curve. Higher orders of n of 7, 9, etc. overlap each other and produce The n equals 5 order corresponds to reinforcement also and is also plotted in mixed colors. Partial overlap of n of 3 and 5 also produce different hues in the uppes order.

rate changes between spheres of different diameters. Standard deviation of 101 readings temperature obtained in 101 readings by several observers of the paint colors around spheres, using theory for rate comparision, and square root of the radius Also plotted in the color curve is the experimental heat transfer and surface from equation I is plus or minus 8 percent.

The Paint is believed to have the following heat transfer characteristics:

K= heat conductions = 10-3+10 tal M=1 (index of restraction)
C= Specific heat = .37 Cal
9m ox 1.2 gm/2 3

there is very little temperature change across the active film, and the film is behaving are between 400 and 1600°K, and that the active film is isothermal in nature, i.e. Use of these characteristics indicates that the wall temperatures for the tests

like a thin film thermocouple. it is important that the angle of incidence be kept constant during various rdg.

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